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Computations involving differential operators and their actions on functions*

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Abstract

This extended abstract further develops the algorithms in [8] and [9] for rewriting expressions involving differential operators. The differential operators that we have in mind arise in the local analysis of nonlinear dynamical systems. In this work, we extend these algorithms in two different directions: We generalize the algorithms so that they apply to differential operators on groups and we develop the data structures and algorithms to compute symbolically the action of differential operators on functions. Both of these generalizations are needed for applications.

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Using Trees To Compute Approximate Solutions to Ordinary Differential Equations Exactly

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Abstract

In this paper, we review some recent work relating families of trees to symbolic algorithms for the exact computation of series which approximate solutions of ordinary differential equations. It turns out that the vector space whose basis is the set of finite, rooted trees carries a natural multiplication related to the composition of differential operators, making the space of trees an algebra. This algebraic structure can be exploited to yield a variety of algorithms for manipulating vector fields and the series and algebras they generate.

1 Introduction

In this paper, we review some recent work relating families of trees to symbolic algorithms for the exact computation of series which approximate solutions of ordinary differential equations. It turns out that the vector space whose basis is the set of finite, rooted trees carries a natural multiplication related to the composition of differential operators, making the space of trees an algebra. This algebraic structure can be exploited to yield a variety of algorithms for manipulating vector fields and the series and algebras they generate.

In Section 3, we introduce and explore the algebraic structure of trees. Section 4 describes a simplification algorithm for the rewriting of symbolic

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expressions involving vector fields. Section 5 describes an algorithm for generating explicitly integrable flows associated with nilpotent Lie algebras. Section 6 exploits the relation between Taylor series and trees to study a class of intrinsic numerical integrators. We begin in Section 2 with some background.

The results surveyed here are the work of a variety of mathematicians: I especially want to mention the contributions of my collaborators Peter Crouch, Matthew Grayson and Richard Larson. The work on algebras of trees and its applications to symbolic computation is joint work with Richard Larson. All of the algorithms described here rest upon this foundation. The work on explicitly integrable flows and nilpotent Lie algebras is joint work with Matthew Grayson. The work on numerical algorithms evolving on groups is joint work with Peter Crouch.

2 Background

Consider a differential equation

$$\dot{x}(t) = E_1(x(t)) + u E_2(x(t)), \qquad x(0) = x^0 \in \mathbf{R}^N,$$
 (1)

where E_1 and E_2 are vector fields and u is a parameter. In applications, u will be either a small perturbation $u = \epsilon$, a control $t \to u(t)$, or simply the constant $u \equiv 1$. Unless the vector fields E_1 and E_2 are very special, no algorithm is known which will return the general solution to the system in closed form. Our objective is to find efficient algorithms to compute various approximate solutions of the differential equation exactly using symbolic computation.

Although the impact of symbolic computation in this area is recent, the connection between the existence of closed form solutions and the approximation of general solutions is a traditional theme, dating back to at least the the nineteenth century. One can distinguish two approaches. One, championed by Lie, is based upon algebra and geometry and concerns us here; the other, championed by Weirstrass and Poincaré, is based upon complex function theory.

Consider a group of transformations acting on \mathbb{R}^N of the form

$$\Phi_s: x_{\mu} = f_{\mu}(x_1, \dots, x_N; s_1, \dots, s_r), \qquad \mu = 1, \dots, N,$$

with the property that the group permutes the solutions of the nonlinear system (1). Lie asked the question [51] and [52], How can information about

the transformation group be used to help integrate the differential equation? To answer this question, Lie introduced the infinitesimal generators of the group

$$A_k(x) = \sum_{\mu=1}^{N} \frac{\partial f_{\mu}}{\partial s_k} \frac{\partial}{\partial x_{\mu}}, \quad 1 \le k \le r$$

and showed that the A_k satisfy

$$[A_i, A_j] = \sum_{k=1}^r c_{ij}^k A_k,$$

where [.,.] is the commutator, or Lie bracket,

$$[A_i, A_j] = A_i A_j - A_j A_i,$$

and the c_{ij}^k are constants. For example, Lie showed that if there is a one parameter group of transformations permuting the solutions of a nonlinear system in the plane (x,t), then the integrating factor for the equation may be read off from the infinitesimal generator.

Since Lie's time, this basic question has contributed to the development of a number of different fields:

- The vector fields A, generate a filtered Lie algebra, which is usually infinite dimensional, and is the infinitesimal version of the continuous pseudogroup of transformations generated by the Φ_s. Prior work has focused on the geometry and structure theory of these algebras; important contributions have been made by Guillemin and Sternberg [37] and [36], and Hermann [52], [53], building upon the earlier work of Cartan, Ehresmann and Spencer.
- Formal sums of iterated powers of vector fields, or Lie series, have been developed by Gröbner [25], [26] and Knapp and Wanner [46], [47] into an operational calculus and used to approximate the solutions of differential equations. Lie transform methods have also been used in perturbation theory by Rand [59] and Meyer [54], in celestial mechanics by Deprit [18], and in particle physics by Dragt [20].
- Explicit series computations of solutions of differential equations have a number of interesting connections with combinatorics. Chen [13] makes uses of the shuffle product, Joyal [45], Labelle [49], Leroux [50], and Viennot [71] employ trees and species, while Rota, Krahaner and Odlyzko [62] exploit the umbral calculus.

- Ritt and Kolchin, building upon earlier work of Picard and Vessiot, developed the field of differential Galois theory. The goal is to obtain a theory describing the solvability of differential equations analogous to Galois' theory describing the solvability of algebraic equations. The survey by Singer [67] provides a good description of this field from the viewpoint of symbolic computation. Other relevant contributions include the beginnings of a differential Gröbner theory [3], [38], analogous to the Gröbner theory in commutative algebra; and a Hopf algebraic interpretation of Picard-Vessiot theory by Takeuchi [70].
- There is now a resurgence of interest in using symmetry groups to help integrate differential equations. This direction of research has been active in the Soviet Union for some time, but during the past decade there has been increased interest in the United States. Important contributions have been made by Olver [57], Schwarz [64], and Bluman and Cole [5]. Closely related is the study by Caviness [66] of conservation laws for differential equations.

During the past several years, there has been increasing interest in symbolic computation and differential equations. Work has proceeded in a number of directions, and is based upon both the Lie and the Weirstrass and Poincaré traditions:

- Zippel [73] is writing a modular symbolic computation system which supports the ability to call high quality numerical routines. Using such a system, he has shown how symbolic algorithms can be used to select appropriate numerical algorithms.
- Guckenhemier [35] has produced the program kaos which allows the user to access a variety of algorithms to investigate a differential equation from the point of view of modern dynamical systems.
- There are a number of programs to compute symmetry groups of differential equations, including one written by Char [12], and the programs SODE and SPDE written by Schwarz [64].
- Abelson and Sussman and their group at MIT [1], [2] have combined techniques in artificial intelligence to produce software which automatically analyzes the qualitative features of a differential equation.
- Wang [72] and Steinberg [68] have developed systems which use symbolic computation to produce high quality, optimized Fortran code to solve differential equations.

$$E_1 \mapsto \bigcup_{E_1}^{C}, \qquad E_2 \mapsto \bigcup_{E_2}^{C}, \qquad E_3 \mapsto \bigcup_{E_3}^{C}$$

Figure 1: The trees associated with the vector fields E_j .

Della-Dora and Tournier [17] have used the fundamental ideas of Ramis
 [58] to produce a system to analyze linear ordinary differential equations. They are now turning their attention to nonlinear systems.

Point of view. The point of view taken here is to focus on the algorithmic aspects of the computation of the vector fields A_j and their brackets and to use this information to develop appropriate algorithms which use exact symbolic techniques to approximately integrate the trajectories of the differential equation. As will become clear, there are a number of interesting points of contact between this approach and the approaches just described.

In the following sections, we review data structures and algorithms for the symbolic computation of the flows of vector fields, and for the symbolic approximation of general flows by flows which can be studied symbolically.

3 Vector fields and the algebra of Cayley trees

In this section, we describe a data structure which is central to the algorithms we give for the symbolic computation of series which approximate the solutions of differential equations. The basic idea is to assign trees to vector fields as illustrated in Figure 1, and then to impose a multiplication on trees which is compatible with the composition of vector fields.

Consider three vector fields

$$E_1 = a_1 D_1 + \cdots + a_N D_N,$$
 $E_2 = b_1 D_1 + \cdots + b_N D_N,$ $E_3 = c_1 D_1 + \cdots + c_N D_N$

where $D_i = \partial/\partial x_i$, and a_i , b_i and c_i are smooth functions on \mathbf{R}^N . Now

$$E_2 \cdot E_1 = \sum b_j(D_j a_i) D_i + \sum b_j a_i D_j D_i$$

and $E_3 \cdot E_2 \cdot E_1$ is equal to

$$\sum c_k(D_kb_j)(D_ja_i)D_i + \sum c_kb_j(D_kD_ja_i)D_i + \sum c_kb_j(D_ja_i)D_kD_i$$

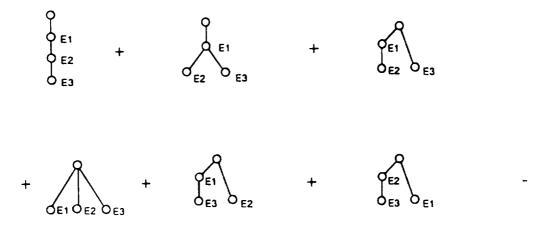


Figure 2: The trees associated with Equation 2.

$$+\sum c_k b_j a_i D_k D_j D_i + \sum c_k b_j (D_k a_i) D_j D_i + \sum c_k (D_k b_j) a_i D_j D_i. \tag{2}$$

Here the sum is for i, j, k = 1, ..., N and hence involves $O(N^3)$ differentiations. It is convenient to keep track of the terms that arise in this way using labeled trees: we indicate in Figure 2 the trees that are associated with the six sums in this expression.

An expression such as

$$[E_{3}, [E_{2}, E_{1}]] = E_{3}E_{2}E_{1} - E_{3}E_{1}E_{2} - E_{2}E_{1}E_{3} + E_{1}E_{2}E_{3}$$
(3)

gives rise in this fashion to 24 trees corresponding to the $24N^3$ differentiations that a naive computation of this expression requires. On the other hand, 18 of the trees cancel, saving us from computing $18N^3$ terms. We are left with $6N^3$ terms of the form $(junk)D_{\mu_1}$. A careful examination of this correspondence between labeled trees and expressions involving the E_j 's shows that the composition of the vector fields E_j 's, viewed as first order differential operators, corresponds to a multiplication on trees. This multiplication is illustrated in Figure 3. It turns out that this construction yields an algebra, which we call the algebra of Cayley trees.

Here is a more precise description for the specialist, following [27] and [30]. Let k denote a field of characteristic 0. We say that a finite rooted tree is labeled with $\{E_1, \ldots, E_M\}$ in case each node, except for the root, is assigned a element from this set. Let $\mathcal{L}T(E_1, \ldots, E_M)$ denote the set of labeled trees and let $k\{\mathcal{L}T(E_1, \ldots, E_M)\}$ denote the vector space whose basis consists of labeled trees in $\mathcal{L}T(E_1, \ldots, E_M)$. Suppose that t_1 ,

$$\bigcup_{E_2} \circ \bigcup_{E_1} = \bigcup_{E_2} + \bigcup_{E_1} \bigcup_{E_2} \bigcup_{E_3} \bigcup_{E_4} \bigcup_{E_4} \bigcup_{E_5} \bigcup_{E_5$$

Figure 3: An example of multiplying two trees.

 $t_2 \in \mathcal{LT}(E_1, \ldots, E_M)$ are trees. Let s_1, \ldots, s_r be the children of the root of t_1 . If t_2 has n+1 nodes (counting the root), there are $(n+1)^r$ ways to attach the r subtrees of t_1 which have s_1, \ldots, s_r as roots to the tree t_2 by making each s_i the child of some node of t_2 . The product t_1t_2 is defined to be the sum of these $(n+1)^r$ trees. It can be shown that this product is associative, and that the trivial tree consisting only of the root is a right and left unit for this product. In [27], we define a comultiplication on $k\{\mathcal{LT}(E_1, \ldots, E_M)\}$ and show that

Theorem 3.1 The vector space $k\{\mathcal{L}T(E_1, \ldots, E_M)\}$ with basis all equivalence classes of finite rooted trees is a cocommutative graded connected Hopf algebra.

We call this algebra the algebra of Cayley trees generated by the set of labeled trees $\mathcal{L}T(E_1,\ldots,E_M)$. The relation between trees and differential operators goes back to Cayley [8], [9]. The coalgebra structure on this space is very similar to the coalgebra structure defined by Joni and Rota [44]. However, the coalgebra structure defined there was defined for individual combinatorial objects, rather than for a class of objects such as the family of rooted trees. Butcher [6] and [7] has also defined a multiplication on the vector space which is dual to the space of trees. This multiplication is closely related to the one just defined.

4 Symbolic evaluation of vector field expressions

When expressions involving vector fields, such as Lie brackets, are written out in coordinates, there is typically a lot of cancellation. Similar cancellation occurs in expressions involving Poisson brackets and when flows are concatenated, as in Campbell-Baker-Hausdorf expansions. In this section, we use the algebra of Cayley trees to exploit this cancellation in order to compute more efficiently formal expressions involving vector fields.

This is the set up. Let R denote a ring of functions. In applications, R is usually either the ring of polynomial functions, rational functions, or C^{∞}

functions. Fix several first order differential operators with coefficients from ${\cal R}$

$$E_{j} = \sum_{\mu=1}^{N} a_{j}^{\mu} D_{\mu}, \quad a_{j}^{\mu} \in R, \quad j = 1, \dots, M$$
 (4)

that are defined in terms of a basis of first order differential operators

$$D_{\mu} = \frac{\partial}{\partial \boldsymbol{x}_{\mu}}, \quad \mu = 1, \ldots, N.$$

Simplifying any expression in the E_j 's using Equation (4) yields a differential operator, which we can view as an element of End R.

We now define a homomorphism from the algebra of expressions in the E, is to the algebra of trees. Indeed, the assignment in Figure 1 extends to an algebra homomorphism from the free associate algebra in the symbols E, to the algebra of Cayley trees $k\{\mathcal{L}T(E_1,\ldots,E_M)\}$. It is straightforward to define a downward-pointing arrow so that the following diagram commutes:

Algorithm 4.1 To rewrite expressions in the first order differential operators E_j in terms of the basis

$$\frac{\partial}{\partial x_{...}}, \quad \frac{\partial^2}{\partial x_{...}, \partial x_{...}}, \dots, \quad \mu_1, \mu_2, \dots = 1, \dots, N,$$

compute the composition of the rightward and downward pointing arrows in the diagram above.

In [28], [29] and [33], we show that the algorithm is much more efficient than naive substitution, which corresponds to computing the diagonal arrow directly. In some common cases, the improvement in efficiency is exponential. We have implemented this and related algorithms in Maple, Mathematica and Snobol.

We illustrate this algorithm by working the example of the last section following [31]: consider a higher order derivation of the form

$$p = E_3 E_2 E_1 - E_3 E_1 E_2 - E_2 E_1 E_3 + E_1 E_2 E_3.$$

Naive simplification requires computing $24N^3$ terms of the form described in Table 1. The image of p in the algebra of Cayley trees contains 24

No. of terms	Form of terms
$8N^3$	coeff. D_{μ_1}
$12N^{3}$	coeff. $D_{\mu_2}D_{\mu_1}$
$4N^3$	coeff. $D_{\mu_3}D_{\mu_2}D_{\mu_1}$

Table 1: Naive computation of the differential operator corresponding to p.

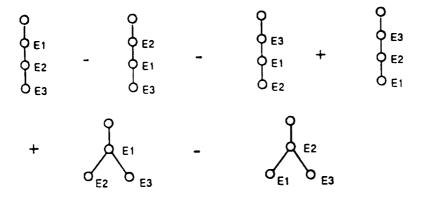


Figure 4: The surviving labeled trees.

trees, six for each of the four terms of p. For example, the six labeled trees corresponding to the first term are given in Figure 2. Eighteen of these trees cancel, leaving the six trees in Figure 4. The corresponding differential operator is equal to

$$\begin{split} &\sum a_{3}^{\mu_{3}} \big(D_{\mu_{3}} a_{2}^{\mu_{2}} \big) \big(D_{\mu_{2}} a_{1}^{\mu_{1}} \big) D_{\mu_{1}} - \sum a_{3}^{\mu_{3}} \big(D_{\mu_{3}} a_{1}^{\mu_{2}} \big) \big(D_{\mu_{2}} a_{2}^{\mu_{1}} \big) D_{\mu_{1}} \\ &- \sum a_{2}^{\mu_{3}} \big(D_{\mu_{3}} a_{1}^{\mu_{2}} \big) \big(D_{\mu_{2}} a_{3}^{\mu_{1}} \big) D_{\mu_{1}} + \sum a_{1}^{\mu_{3}} \big(D_{\mu_{3}} a_{2}^{\mu_{2}} \big) \big(D_{\mu_{2}} a_{3}^{\mu_{1}} \big) D_{\mu_{1}} \\ &+ \sum a_{3}^{\mu_{3}} a_{2}^{\mu_{2}} \big(D_{\mu_{3}} D_{\mu_{2}} a_{1}^{\mu_{1}} \big) D_{\mu_{1}} - \sum a_{3}^{\mu_{3}} a_{1}^{\mu_{2}} \big(D_{\mu_{3}} D_{\mu_{2}} a_{2}^{\mu_{1}} \big) D_{\mu_{1}}, \end{split}$$

and contains $18N^3$ fewer terms of the form indicated in Table 2 than does the naive computation of p. An example of the cancellation of labeled trees is given in Figure 5.

To summarize: we have defined an algebraic structure on families of trees which mirrors the algebraic structure of formal expressions in the variables E_j , but which alleviates the need for computing intermediate expressions which cancel when the noncommuting E_j 's are expressed in terms of the

$$- \bigcup_{E_1}^{E_2} + \bigcup_{E_3}^{E_2} = 0$$

Figure 5: The term $E_3E_1E_2$ contributes the first labeled tree and the term $E_1E_2E_3$ contributes the second, which cancel.

No. of terms	Form of terms
$2N^3$	coeff. D_{μ_1}
12N ³	coeff. $D_{\mu_2}D_{\mu_1}$
4 N ³	coeff. $D_{\mu_1} D_{\mu_2} D_{\mu_1}$

Table 2: Terms in the computation p which cancel.

commuting D_{μ} 's. In the following sections, we will see other expressions of this simple idea.

Algorithm 4.1 can be extended in several different directions.

1. The examples above concern vector fields defined on \mathbb{R}^N . It is possible to work out similar algorithms for vector fields defined on more general objects, such as Lie groups. This is important for applications in robotics and rigid body dynamics. For example, the group G = SO(3) is the appropriate configuration space for a rotating rigid body. To be more specific, assume the vector fields are of the form

$$E_j = \sum_{\mu=1}^N a_j^{\mu} Y_{\mu},$$

where the Y_{μ} are left-invariant vector fields on G and the a_{j}^{μ} are smooth functions on the group. In this case, the Cayley algebras are generated by *ordered*, labeled trees [16]. Roughly speaking, the trees are ordered since the vector fields Y_{μ} no longer commute.

2. The natural action of differential operators on functions turns the ring of functions R into a module. In this same way, the trees have a natural action on R, as indicated in the Diagram 5. It turns out [34] that this gives R the structure of K/k-bialgebra, as introduced by Nichols [55]

- and [56]. These types of algebras are closely related to differential algebras.
- 3. It is a basic fact that the local properties of the nonlinear system (1) are determined by the algebraic properties of the higher order iterated Lie brackets; see, for example, [40]. Unfortunately, due to intermediate expansion swell, it is often difficult to compute these using current computer algebra systems. It turns out that higher order Lie brackets not only involve the cancellation of all terms above the first order but also the cancellation of some of the first order terms. Algorithm 4.1 can be used so that the terms arising in these first order cancellations need not be computed.

5 Exponentials, Lie brackets, and nil flows

Some differential equations have the property that their flows can be integrated symbolically in closed form. For example, this holds for differential equations of the form (1), if E_1 and E_2 are homogeneous in the appropriate sense and generate a graded, nilpotent Lie algebra. In this section, we give an algorithm which, given an appropriate nilpotent Lie algebra, returns vector fields on \mathbb{R}^N with polynomial coefficients which generate the Lie algebra. This leads to an interesting class of differential equations whose flows can be integrated symbolically in closed form. At the end of the section, we look at several applications of this algorithm.

By the third fundamental theorem of Lie [63], we know that a nilpotent Lie algebra arises from some Lie algebra of vector fields. What is not obvious is how to construct such vector fields. Nilpotent Lie algebras of vector fields have been used as an important tool in partial differential equations by Folland and Stein [21], Rothschild and Stein [60], and Rockland [61]; and in control theory by Krener [48], Hermes [41], [42], and Crouch [15], [14]. We will see below how they are also a useful tool in developing symbolic-numeric algorithms to integrate flows.

Our goal is to describe a natural representation of nilpotent Lie algebras on vector fields on Euclidean space with polynomial coefficients. To define this representation, we define a basis of Hall trees on generators E_1 and E_2 recursively as follows:

1. basis elements consist of rooted, binary trees, with all nodes, except the root, labeled with E_1 , E_2 , E_3 , ... satisfying

- (a) all right children are leaves and labeled with either E_1 or E_2
- (b) the sequence formed by the labels of the right leaves at increasing distance from the root is nonincreasing
- 2. the two rooted trees consisting of a root and a single (left) child labeled E_i , for i = 1, 2 are in the basis and of length 1
- 3. if we have defined basis elements t_i of length $1, \ldots, r-1$, they are simply ordered so that $t_i < t_j$ in case the length of t_i is less than the length of t_j .
- 4. a tree consisting of a root with a single (left) child labeled E_i is in the basis provided that E_i 's left child is in the basis and of lower order.

To each such tree corresponds an element E_i in the Hall basis [39] of the free, nilpotent Lie algebra on two generators. Figure 6 illustrates how the basis of Hall trees leads naturally to a representation of the free nilpotent Lie algebra generated by E_1 and E_2 on the space of vector fields on \mathbf{R}^N with polynomial coefficients. See [22] for further details of the following algorithm:

Algorithm 5.1 Fix r > 1 and say that the free, nilpotent Lie algebra of two generators of rank r has dimension N. Let E_1, \ldots, E_N denote the Hall basis. Then the vector fields on \mathbf{R}^N defined as in Figure 6 satisfy

- 1. the Lie algebra they generate is isomorphic to the free, nilpotent Lie algebra on two generators of rank r
- 2. any trajectory of the nonlinear system

$$\dot{x}(t) = E_1(x(t)) + u(t) E_2(x(t)), \qquad x(0) = x^0 \in \mathbb{R}^N,$$

can be computed in closed form in terms of quadratures of the function $t \to u(t)$

3. there exist constants $\alpha_1, \ldots, \alpha_N$ such that

$$\exp(E_2)\exp(E_1) = \exp\left(\sum_{i=1}^N \alpha_i E_i\right),$$

and the α_i can be computed by solving a lower triangular linear system.

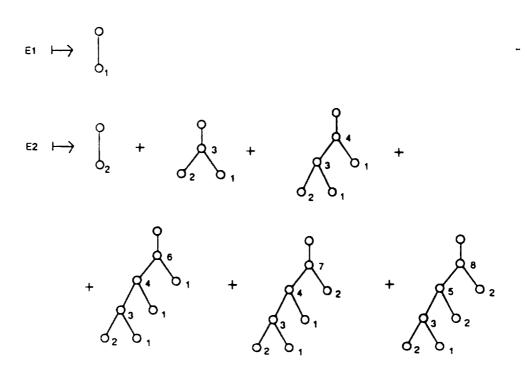


Figure 6: The vector fields arising from the basis of Hall trees. The first tree is sent to D_1 , while the sum of the trees is sent to $D_2 - x_1 D_3 + \frac{1}{2} x_1^2 D_4 + x_1 x_2 D_5 - \frac{1}{6} x_1^3 D_6 - \frac{1}{2} x_1^2 x_2 D_7 - \frac{1}{2} x_1 x_2^2 D_8$.

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As a note, this algorithm was found only after several months of experimentation with Maple and required the careful study of the free, nilpotent Lie algebra on two generators of rank 5, which is 23 dimensional.

We conclude this section with some remarks describing several applications of this algorithm and related work.

- 1. Locally approximating a nonlinear system by an explicitly integrable nilpotent one yields a number of integration algorithms, which we refer to as piecewise nilpotent integration algorithms (PWNI). The basic idea [23], [24] is to approximate locally a nonlinear system at a given point by an explicitly integrable nilpotent one in which the computations can be done symbolically in closed form, and to patch together the various nilpotent approximations at nearby points using a standard numerical algorithm. Preliminary work indicates that this leads to symbolic-numeric algorithms for the path planning problem and efficient algorithms to integrate neighborhoods of trajectories around a given fixed reference trajectory.
- 2. Algorithm 5.1 also provides an efficient means for computing the concatenation of flows. Write $x(t) = \exp E \cdot x^0$ for the flow of the nonlinear system

$$\dot{x}(t) = E(x(t)), \quad x(0) = x^0,$$
 (6)

The Campbell-Baker-Hausdorff formula expresses the concatenation of two flows as a single flow:

$$\exp(tE_2)\exp(tE_1) = \exp(tE_2 + tE_1 + 1/2t^2[E_2, E_1] + 1/12[[E_2, E_1], E_1] - 1/12[[E_2, E_1], E_2] + \cdots). \tag{7}$$

Consider the equation

$$\exp(E_2)\exp(E_1) = \sum_{i=1}^N c_i E_i,$$

where E_i are the vector fields produced by the algorithm. Since all flows of the vector fields E_1 and E_2 are explicitly integrable in closed form, this reduces the computation of the c_i to the solution of a linear system, which is lower triangular.

3. We can also use Algorithm 5.1 to derive a class of numerical integrators, which are sometimes known as splitting methods. For example,

suppose that E_1 and E_2 are separately integrable in closed form, but $E_1 + E_2$ is not. Then using the algorithm, we can compute constants τ_1 such that

$$\exp(\tau_7 E_1) \cdot \exp(\tau_6 E_2) \cdot \exp(\tau_5 E_1) \cdot \exp(\tau_4 E_2)$$

$$\exp(\tau_3 E_1) \cdot \exp(\tau_2 E_2) \cdot \exp(\tau_1 E_1) \cdot = \exp(E_1 + E_2) + O(t^5).$$
 (8)

This formula yields a numerical integrator for our original nonlinear system (1) (with $u \equiv 1$). This algorithm is used in accelerator physics [65].

4. Recently, Strichartz [69] has shown that the solution of the initial value problem

$$\dot{x}(t) = E(t, x(t)), \quad x(0) = x^{0}$$

can be written as

$$x(t) = \exp(G(t))x^0,$$

with

$$G(t) \sim \sum_{r=1}^{\infty} \sum_{\sigma} c_{r,\sigma}$$

$$\int \left[\cdots \left[E(s_{\sigma(1)}, E(s_{\sigma(2)}) \cdots\right] E(s_{\sigma(r)})\right] ds,$$

and where σ ranges over the symmetric group on r symbols, the integration region is a simplex in \mathbf{R}^r , E(s) denotes the vector field $E(s, \cdot)$, and exp is defined as in Equation 6. Formulas of this type date back to Chen [13] and appear to be related to Algorithm 5.1.

5. F. Bergeron, N. Bergeron, and A. M. Garsia have also exploited a relation between trees and polynomials in their study of free Lie algebras; see [4] and the references cited there.

6 Taylor series and intrinsic integrators

Although nonlinear systems often conserve quantities such as energy or angular momentum, most numerical integrators do not. Similarly, nonlinear systems typically evolve on some underlying geometric space, such as a Lie group or homogeneous space, but most numerical integrators do not remain in such a space.

Recently, there has been a flurry of activity related to numerical integrators preserving the symplectic structure, sparked off by the work of

Channell [10] and Scovel [11], and based upon earlier work of deVogeleaere [19]. These types of numerical schemes have found important applications in the long term study of orbits in accelerator physics and in other areas [65]. The derivation of these integrators typically involves the symbolic computation of Taylor series and generating series for the symplectic transformation which is the update for the numerical integrator.

Consider a numerical integrator for a differential equation

$$\dot{x}(t) = E(x(t)), \qquad x(0) = p \in M$$

evolving on a space M. Call a numerical integrator intrinsic in case $x_n \in M$ implies $x_{n+1} \in M$, for $n \ge 1$, where x_n is the approximation to the trajectory x(t) at time t_n . One means of deriving intrinsic numerical integrators is to mimic the derivation of standard numerical integrators, but to impose additional constraints on the scheme to satisfy the added condition that the points x_n remain in the space M. This typically involves the careful study of the Taylor series of the solution.

This can be done by using the Cayley algebra of trees, as briefly indicated in [32]. As an illustration of this, we consider intrinsic Runge-Kutta type algorithms evolving on a Lie group G, following [16]. Let g denote its Lie algebra, and let Y_1, \ldots, Y_N denote a basis of g. We give an algorithm to approximate solutions to differential equations evolving on G of the form:

$$\dot{x}(t) = E(x(t)), \qquad x(0) = p \in G,$$

where

$$E=\sum_{\mu=1}^N a^{\mu}Y_{\mu},$$

and the $=a^{\mu}$ are analytic functions on G. Let $\exp(tE) \cdot x$ denote the solution x(t) at time t. The algorithms depends upon constants c_i and c_{ij} , for $i=1,\ldots,k$ and j< i. For fixed constants, define the following elements of the Lie algebra g

$$\bar{E}_{1} = \sum_{\mu=1}^{N} a^{\mu}(x_{n}) Y_{\mu} \in g$$

$$\bar{E}_{2} = \sum_{\mu=1}^{N} a^{\mu}(\exp(hc_{21}\bar{E}_{1}) \cdot x_{n}) Y_{\mu} \in g$$

$$\bar{E}_{3} = \sum_{\mu=1}^{N} a^{\mu}(\exp(hc_{32}\bar{E}_{2}) \cdot \exp(hc_{31}\bar{E}_{1}) \cdot x_{n}) Y_{\mu} \in g \dots$$



Figure 7: Trees associated with third order terms in a Taylor series.

These arise by "freezing the coefficients" of E at various points along the flow of E

Algorithm 6.1 Given an initial point x0 on the group, define

$$x_{n+1} = \exp hc_k \tilde{E_k} \cdots \exp hc_1 \tilde{E_1} x_n,$$

for $n \geq 0$.

Notice that if we assume the exponential $\exp(h\bar{E_i})$ maps the Lie algebra to the Lie group exactly, then this algorithm is intrinsic. For a group such as G = SO(3), there are classical closed form expressions for the exponential map and Algorithm 6.1 yields an intrinsic integrator. Notice also that if G is the abelian group RN, then the algorithm becomes the classical Runge-Kutta algorithm.

The first step is to derive the equations that the coefficients c_i and c_{ij} must satisfy in order for the algorithm to yield an rth order numerical integrator. This can easily be done using the Cayley algebra of ordered trees [16]. The trees are ordered since the vector fields Y_{μ} do not commute. Assume for the moment that $G = \mathbb{R}^{N}$ and consider the terms in the

Taylor series

$$x(t+h) - x(t) = h\dot{x}(t) + \frac{h^2}{2!}\ddot{x}(t) + \frac{h^3}{3!}x^{(iii)}(t) + \cdots$$

$$= hE + \frac{h^2}{2!}DE \cdot E + \frac{h^3}{3!}\left(DE \cdot DE \cdot E + D^2E(E, E)\right) + \cdots$$

Notice that there is a natural correspondence between trees and terms in the series. For example, the $h^3/3!$ terms are associated with the two labeled trees in Figure 7. This observation goes back to at least Cayley [8], [9].

We now generalize this to a Lie group following [16]. Recall that Diagram 5 induces an action of trees on the ring of analytic functions on G. Using this action, we can now state the

Lemma 6.1 Let α denote the tree consisting of a root with a single child labeled F. Then for any analytic function f on the group and for sufficiently small t.

$$f(\exp(tF) \cdot x) = \exp(t\alpha) \cdot f \mid_{x}.$$

Notice that if G is Euclidean space, and if the functions f are the coordinate functions x_1, \ldots, x_n , then this becomes the familiar Taylor series.

Using this lemma, it is now easy to compute the equations that the coefficients c_i and c_{ij} must satisfy. In spirit, this is similar to Butcher's use of trees to analyze higher order Runge-Kutta algorithms in Euclidean space [6], [7].

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